



## Influence of geometric and electronic characteristics of TiO<sub>2</sub> electrodes with nanotubular array on their photocatalytic efficiencies

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### ABSTRACT

The present work compares the efficiencies of titania nanotubes (NT) formed by anodizing of titanium at 20 V for 45 min in three kinds of electrolytes: ethylene glycol–ammonium fluoride–water (NT<sub>1</sub>), tetrabutylammonium fluoride–formamide–water (NT<sub>2</sub>) and sodium sulfate–ammonium fluoride–water (NT<sub>3</sub>). The geometric characteristics were investigated by scanning electron microscopy (SEM). The photoactivities of the three layers were compared through photo-potential, photo-current measurements in an alkaline medium and in the presence of methanol as a hole scavenger. Impedance spectrometry was used for investigation of conduction and charge storage properties of the titania layer. Finally the photocatalytic properties were compared through the degradation rate of an azo dye, such as Amido Black (AB).

In comparison to NT<sub>2</sub> and NT<sub>3</sub>, the NT<sub>1</sub> layer showed not only the highest photopotential, but also the highest photocurrent response. Photocatalytic results showed that the degradation of dye proceeds much faster in the presence of NT<sub>1</sub> as compared with other TiO<sub>2</sub> nanotube samples. Nevertheless photoinduced surface states were found to be more important in NT<sub>2</sub>, in comparison with the two other layers. All these results are discussed taking into account the geometric characteristics of the nanotubular layers, not only the specific surface on one hand, but also the solid fraction of the external surface, which determines the level of light absorption able to generate photo-charges.

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### 1. Introduction

TiO<sub>2</sub> is a widely used photocatalyst for degradation of organic pollutants in waste water, because of its high photosensitivity, long-term stability and non-toxicity [1]. The photocatalytic properties of TiO<sub>2</sub> nanotubular thin layers (NT) have been studied by several groups, and much higher decomposition efficiency than for a Degussa P25 nanostructured layer under similar conditions has been reported [2–4]. Two key factors influence the photocatalytic efficiency: the first one is the huge specific area of the nanotubular layer which allows increased adsorption of organic molecules during the photocatalytic reaction [4]. The second determining factor is the transport rate of the photogenerated charges through the layer. Nanotubular films made of oriented one-dimensional structures aligned perpendicular to the substrate are assumed to show improved charge-collection efficiency by promoting faster charge transport and slower recombination

[5–8]. In nanostructured layers, the interfacial kinetics are also drastically affected by the presence of electron traps in the layer, and the charge transport in such films can be controlled by the mechanism of trapping–detrapping [9–18].

Among the different fabrication methods of NT, the titanium anodization method in fluoride-based baths provides good mechanical adhesion and good electronic conductivity, since it directly grows from the titanium metal substrate [2,19–24].

The thickness and morphology of such TiO<sub>2</sub> film are strongly affected by the electrochemical conditions (anodization voltage, duration of anodization, ...) and the composition of the electrolyte (fluoride concentration, pH, water content in the electrolyte and presence of an organic solvent, ...). Since the as-prepared samples show an amorphous structure, post-treatments are necessary to control the crystallinity of the nanotubes. Thermal annealing in the 450–600 °C temperature range is usually carried out to get a well crystallized anatase structure. Yu and Wang [25] showed that the calcinations temperature exhibits a great influence on the photocatalytic activity. A vapour-thermal treatment at 180 °C instead of calcination [26], or a chemical treatment of the tube walls applied after calcinations [27] has been shown to further improve the photocatalytic activity.

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Comparison between the morphology and the properties of nanotubular TiO<sub>2</sub> layers obtained in different electrolytes has been reported by Zlamal et al. [28]. They suggest that thick tube walls are needed for good catalytic performance since they allow more efficient separation of electron-hole pairs. Smooth tube walls are also beneficial for an improved electron diffusion length in the layer [5]. The presence of a high density of deep localized electronic states in the gap of nanotubular TiO<sub>2</sub> layers has been evidenced in [23], but no clear correlation between the presence of electronic states or traps, and the photocatalytic activity of the layer has been reported.

The aim of the present work is thus to contribute to establish a correlation between the morphology of the NT layers, the density of surface states present in the layers and their photocatalytic activity. Three kinds of NT layers with similar thicknesses were synthesized. The first kind of nanotubes (NT<sub>1</sub>) was grown in ethylene glycol in order to get tubes with smooth walls. The second kind (NT<sub>2</sub>) was obtained in a tetrabutylammonium fluoride solution prepared in formamide in order to get rough walls. The third kind of NT was synthesized in an aqueous medium in the absence of any organic solvent. The photoactivity of the three layers was compared through photo-potential and photo-current measurements. Impedance spectrometry was also used for investigation of conduction and charge accumulation properties. All the electrochemical experiments were performed in an alkaline medium and in the presence of methanol as a hole scavenger. Finally the photocatalytic properties were compared through the degradation rate of an azo dye, such as Amido Black (AB).

## 2. Experimental

### 2.1. Synthesis of TiO<sub>2</sub> nanotubular layers

Anodization was carried out at 20 V for 45 min, at room temperature using a DC power source in a two-electrode cell, where a Pt foil served as the counter electrode and a titanium sheet was used as the working electrode. Prior to anodization, the titanium foil was first mechanically polished with different abrasive papers, sonicated in acetone, ethanol and deionized (DI) water, and then dried in a N<sub>2</sub> stream.

Three kinds of TiO<sub>2</sub> nanotube arrays films were made by anodic oxidation of a titanium rod (2 mm thickness, 15 mm diameter, 99.6% purity) in the following solutions:

- for NT<sub>1</sub>: ethylene glycol containing 3 wt% NH<sub>4</sub>F and 2 vol% H<sub>2</sub>O [29];
- for NT<sub>2</sub>: 0.27 mol/L tetrabutylammonium fluoride solution prepared in formamide containing 3% deionized water [30,31];
- for NT<sub>3</sub>: 1 mol/L Na<sub>2</sub>SO<sub>4</sub> with 1 wt% NH<sub>4</sub>F in water [32].

Finally, as-anodized samples were first thoroughly rinsed with ethanol and with water, dried under nitrogen then annealed in air at 500 °C during 2 h with a ramp-up and ramp-down rate of 1 °C/min.

### 2.2. Crystalline structure of the electrodes

The crystalline structures of the annealed electrodes were studied using an OLYMPUS-JOBIN YVON-Raman spectrometer. The surface morphology of films was observed using a scanning electron microscope SEM-FEG Ultra 55.

### 2.3. Photoelectrochemical procedure

Photoelectrochemical characterization was performed using a VoltaLab 40 PGZ301 potentiostat (Radiometer Analytical), connected with a computer that uses VoltaMaster 4.0 software for data

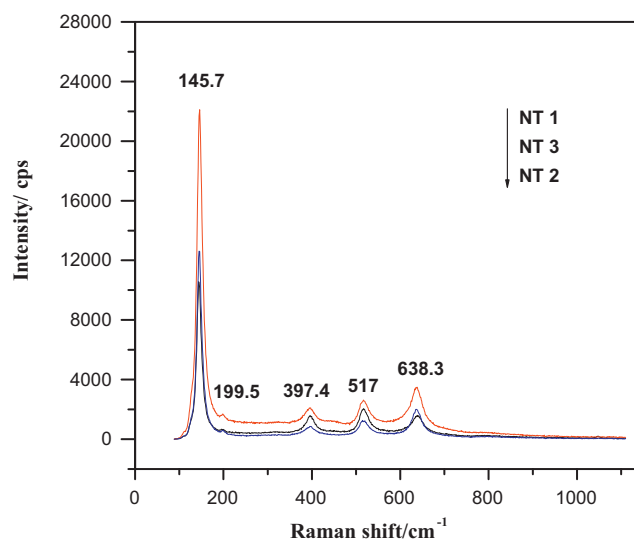


Fig. 1. Raman spectra for the three nanotubular TiO<sub>2</sub> layers.

acquisition. A 100 mL cell made of quartz was used as photoelectrochemical cell with the TiO<sub>2</sub>/Ti electrode as working electrode, a Pt counter-electrode and a saturated calomel reference electrode (SCE). All potentials are quoted versus SCE. The geometric surface of the working electrode was 0.64 cm<sup>2</sup>. The electrolyte was an air-saturated aqueous solution with 0.1 mol/L NaOH, with and without methanol. A 150 W Xenon lamp was employed as excitation source. The electrochemical investigations were performed in the dark and under light excitation in solutions containing different methanol concentrations (0–100 mmol/L). Electrochemical impedance spectroscopy (EIS) measurements were carried out at the open-circuit potential. A sinusoidal AC perturbation of 10 mV was applied to the electrodes over the frequency range of 20,000–0.1 Hz. The experimental data from EIS diagrams were analysed using the software, ZsimpWin3.2.

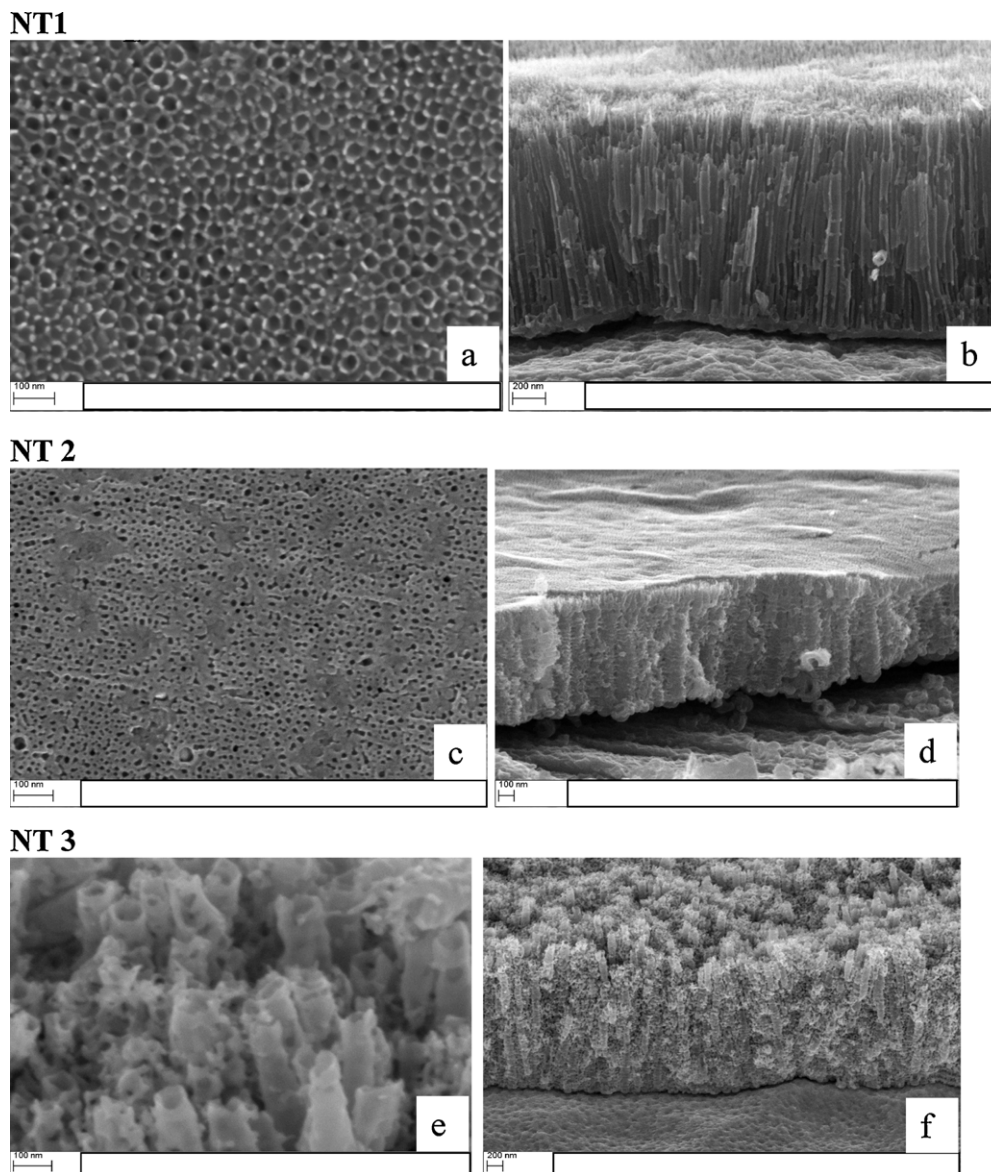
## 3. Results and discussion

### 3.1. Raman spectroscopy analyses

Raman spectroscopy is a technique that can clearly identify the crystal forms of titania (anatase, rutile and brookite) even at very low levels. In fact the Raman spectrum of anatase single crystal shows six Raman bands at 144 (*E<sub>g</sub>*), 197 (*E<sub>g</sub>*), 399 (*B<sub>1g</sub>*), 513 (*A<sub>1g</sub>*), 519 (*B<sub>1g</sub>*) and 639 cm<sup>-1</sup> (*E<sub>g</sub>*) [33]. The rutile phase is characterized by four Raman active modes at 143 (*B<sub>1g</sub>*), 447 (*E<sub>g</sub>*), 612 (*A<sub>1g</sub>*), and 826 cm<sup>-1</sup> (*B<sub>2g</sub>*) [34]. The brookite phase has three Raman bands at 247, 326 and 366 cm<sup>-1</sup> [35]. In this work, the Raman spectra of the three titania nanotubes samples mainly show the typical bands of anatase phase (see Fig. 1).

### 3.2. SEM analysis

Fig. 2 shows a typical SEM image of a TiO<sub>2</sub> nanotube layer prepared by anodization of titanium at 20 V for 45 min in the three different electrolytes. It can be seen from the top-view images that the nanotube array produced using ethylene glycol (NT<sub>1</sub>) is more uniform, with smooth tube walls, well-oriented over the substrate. NT<sub>2</sub> tubes obtained in tetrabutylammonium fluoride and formamide solution show rough walls and the individual tubes are connected with each other via rings at the side walls. NT<sub>3</sub> tubes formed in aqueous medium in the absence of organic solvent are more spaced from each other in comparison with NT<sub>1</sub> and NT<sub>2</sub>. They



**Fig. 2.** SEM images of TiO<sub>2</sub> nanotubes layers formed on Ti substrate: (a, c and e) top-view; (b, d and f) cross-section of NT<sub>1</sub>, NT<sub>2</sub> and NT<sub>3</sub>, respectively.

seem less well-oriented over the substrate and the diameter of the tubes increases from top to the bottom of the layer. It is also clear from these pictures that for the three types of layers, the tubes are open on the top and closed at the bottom of the layer.

### 3.3. Geometric characteristics of the layers

The mean geometric parameters of the tubes: length ( $L$ ), outer ( $D_{out}$ ) and inner ( $D_{in}$ ) diameter, wall thickness ( $w$ ), void thickness ( $y$ ) between the tubes are reported in Table 1. It is noticeable that the wall thickness ( $w$ ) has similar values in the three layers.

According to [30,36], a geometric model with a regular network of identical and equally spaced nanotubes can be used for the estimation of the geometric roughness factor  $H$ , defined as the ratio of the cylindrical inner and outer surfaces relative to the corresponding projected area.

Considering the parameters  $L$ ,  $y$ , and  $R_1 = (1/2)D_{int}$ ;  $R_2 = R_1 + (1/2)w$ , the  $H$  factor is calculated as following:

$$H = \frac{8\pi LR_2}{\sqrt{3}(4R_2 - 2R_1 + y)^2} \quad (1)$$

The calculated geometric factor  $H$  and the corresponding specific surface ( $S_{spec} = H \times \text{geometric surface}$ ) are also reported in Table 1. Though NT<sub>3</sub> shows a higher tube length and a higher void thickness between the tubes, the specific area are of the same order of magnitude in the three layers, with a value slightly larger for NT<sub>1</sub>.

Moreover, using the same geometric model the tube density is given by:

$$d_{NT} = \frac{2}{\sqrt{3}(4R_2 - 2R_1 + y)^2} \quad (2)$$

The solid fraction  $S_{\%}$  is the surface ratio at the top of the layer able to absorb incident photons and is calculated according to Eq. (3):

$$S_{\%} = \frac{\pi}{4}(D_{out}^2 - D_{in}^2)d_{NT} \quad (3)$$

The  $d_{NT}$  and  $S_{\%}$  values are also reported in Table 1.

### 3.4. Open-circuit potential ( $E_{OCP}$ )

Fig. 3(a) shows the evolution of the open-circuit potential with the concentration of methanol, in a 0.1 mol/L air-saturated NaOH

**Table 1**  
Geometric parameters of the 3 layers with nanotube arrays (geometric surface = 0.64 cm<sup>2</sup>). *L*: tube length; *D*<sub>in</sub>: inner diameter; *D*<sub>out</sub>: outer diameter; *y*: void thickness; *w*: wall thickness.

Electrode	<i>L</i> (nm)	<i>D</i> <sub>in</sub> (nm)	<i>D</i> <sub>out</sub> (nm)	<i>y</i> (nm)	<i>w</i> (nm)	<i>H</i>	Specific surface (cm <sup>2</sup> )	<i>d</i> <sub>NT</sub> cm <sup>-2</sup>	<i>S</i> <sub>%</sub>
NT <sub>1</sub>	1250	40	70	20	15	61	39	1.4 × 10 <sup>10</sup>	0.36
NT <sub>2</sub>	800	23	63	20	20	36	23	1.7 × 10 <sup>10</sup>	0.47
NT <sub>3</sub>	1700	70	106	50	18	45	29	0.47 × 10 <sup>10</sup>	0.23

solution, in the dark and under illumination. The values were taken when the steady state was reached, after 30 min of immersion.

In the absence of UV light, the OCP value is slightly more negative for NT<sub>1</sub> (−0.55 V) than for NT<sub>2</sub> (−0.50 V) and NT<sub>3</sub> (−0.52 V). In the three cases, it is independent of the concentration of methanol. On the contrary, under illumination, a clear shift of the OCP towards more negative values is observed, when the concentration of methanol, is increased, reaching a value of −1.05 V for NT<sub>1</sub> in 100 mmol/L of methanol. Fig. 3(b) shows the variation of photopotential (difference between OCP in the dark and OCP under lightening) as a function of methanol concentration. It is noticeable that NT<sub>1</sub> shows the highest photo-induced potential shift among the three nanolayers.

Under open-circuit conditions, the photoanodic current must be associated with a cathodic reaction draining the photogenerated electrons. The increase of the photopotential with increasing methanol concentration can thus be attributed to the increase of the anodic reaction (hole consumption by methanol), when methanol concentration increases. Simultaneously, more electrons are available to accumulate in the electrode. Comparing the photopotential

values in the three types of layers, we must admit that charge separation between photogenerated holes and electrons is more efficient in NT<sub>1</sub> layers, in comparison with NT<sub>2</sub> and NT<sub>3</sub>.

### 3.5. Photocurrent measurements

Photocurrent measurements were carried out at 0.8 V, in a potential range where saturation of the photocurrent was reached and charge recombination in the semiconductor could be neglected. In these conditions the photocurrent *I*<sub>photo</sub> can be estimated from empirical relation (4).

$$I_{\text{photo}} = K S_{\text{spec}} S_{\%} L \alpha \quad (4)$$

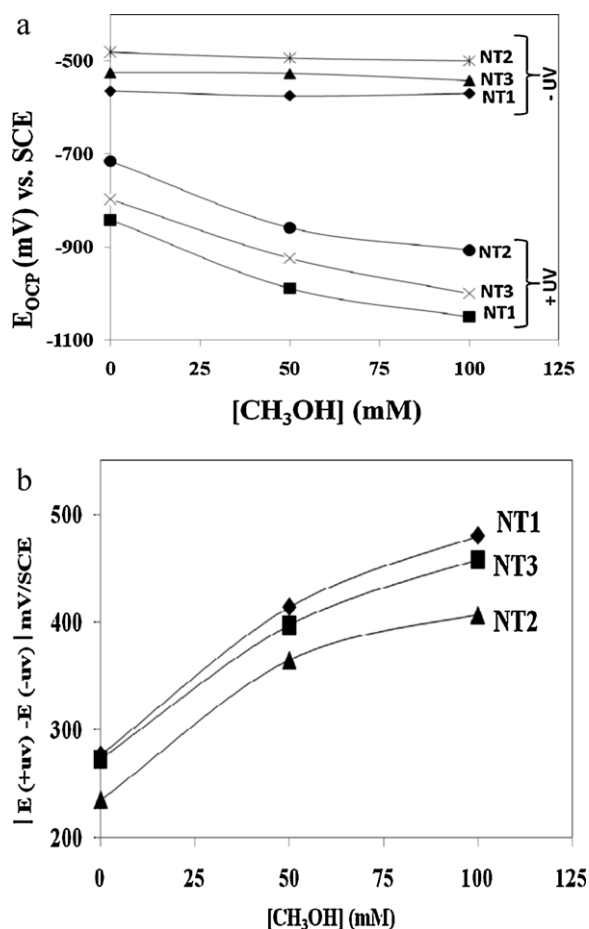
where  $\alpha$  is the effective absorption coefficient for the spectral range of illumination,  $\alpha = 0.5 \mu\text{m}^{-1}$  in anatase [5]. *S*<sub>spec</sub>, *L* and *S*<sub>%</sub> values were defined above and are reported in Table 1.

*K* depends on the incident photon flux  $\Phi_0$  and on the coverage rate  $\theta$  of reacting species at the surface of the nanotubes according to Eq. (5):

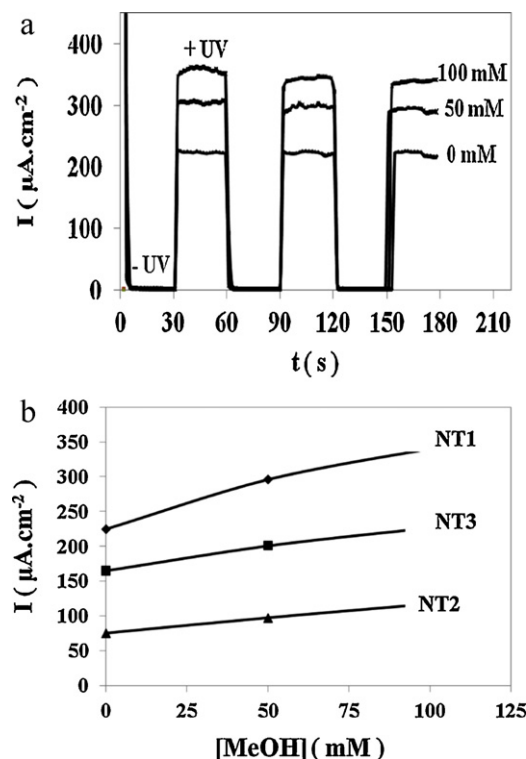
$$K = q\Phi_0\theta \quad (5)$$

It has dimension of a current density (A cm<sup>-2</sup>) with *q*, the elementary charge.

Fig. 4 shows that the photocurrent for NT<sub>1</sub> is always higher than those of NT<sub>2</sub> and NT<sub>3</sub> samples, in the absence and in the presence of methanol. In absence of methanol, the reacting species are mainly adsorbed OH<sup>-</sup>.



**Fig. 3.** (a) Open-circuit potential reached after 30 min in 0.1 mol/L NaOH as a function of methanol concentration, in the dark (−UV) and under illumination (+UV), (b)  $\Delta E$  vs. methanol concentration.



**Fig. 4.** (a) Photocurrent–time characteristics at 0.8 V/SCE for NT<sub>1</sub> sample; (b) photocurrent vs. methanol concentrations for the three nanotubular samples.

**Table 2**

Comparison between experimental and calculated photocurrent values on the bases of geometric characteristics of the NT layers.

Electrode	$I_{\text{photo}}$ ( $\mu\text{A}$ ) experimental	$L\alpha$ ( $\alpha = 5000 \text{ cm}^{-1}$ )	$S_{\text{spec}} S_{\%} L\alpha$	$I_{\text{photo}}$ ( $\mu\text{A}$ ) calculated
NT <sub>1</sub>	230	0.62	8.7	230
NT <sub>2</sub>	70	0.40	4.1	108
NT <sub>3</sub>	160	0.87	5.8	153

Under a constant light intensity, and for a given concentration in reacting species,  $K$  remains constant. According to Eq. (4), in the absence of methanol, the photocurrents for NT<sub>2</sub> and NT<sub>3</sub> were calculated, extracting  $K$  from the experimental photocurrent measured for NT<sub>1</sub>. The experimental and calculated values are compared in Table 2 and it appears that the differences between the measured photocurrent values can be attributed to differences in the geometric characteristics of the layer.

As shown in Fig. 4(a and b), the photocurrent responses also depend on methanol concentration. In Fig. 4(a), reported for NT<sub>1</sub> as an example, the rise and fall of the photocurrent at 0.8 V correspond well to the switch on and off of illumination. The variation of the photocurrent at this applied bias, as a function of methanol concentration is reported in Fig. 4(b), for the three types of electrodes. The observed increase of  $I_{\text{photo}}$  with methanol concentration is in good agreement with the results of Zhang et al. [37] showing a linear dependence of the saturation photocurrent with methanol concentration in the range of 0–100 mmol L<sup>-1</sup>. They suggest that the steady-state photocurrent is first order with respect to methanol concentration in this low concentration range, and is limited by the light intensity at higher concentrations. An increase with methanol concentration of the oxidation rate by photogenerated holes is also in agreement with the shift of the open-circuit potential under lightening when the methanol concentration is increased.

### 3.6. Voltammetry

Under cathodic polarization, in aerated solutions several features can be observed for titanium oxide: besides expected oxygen and water reduction, a cation in-diffusion reaction, combined with reduction of the oxide, occurs simultaneously [38,39]. This later process is often electrochemically reversible and was reported to modify the electronic properties of the layer and to generate additional energy states within the band gap [30].

Fig. 5(a) shows a  $i=f(E)$  curve performed in air-saturated 0.1 M NaOH solution, in the dark, starting from open-circuit potential (about -0.5 V) towards -2 V, followed by a reverse scan in the anodic range. A current plateau due to oxygen reduction is observed between -1 V and -1.5 V. At more negative potentials hydrogen evolution occurs. No specific feature for proton in-diffusion or oxide reduction can be distinguished, thus if they occur the corresponding current is included in the total cathodic current.

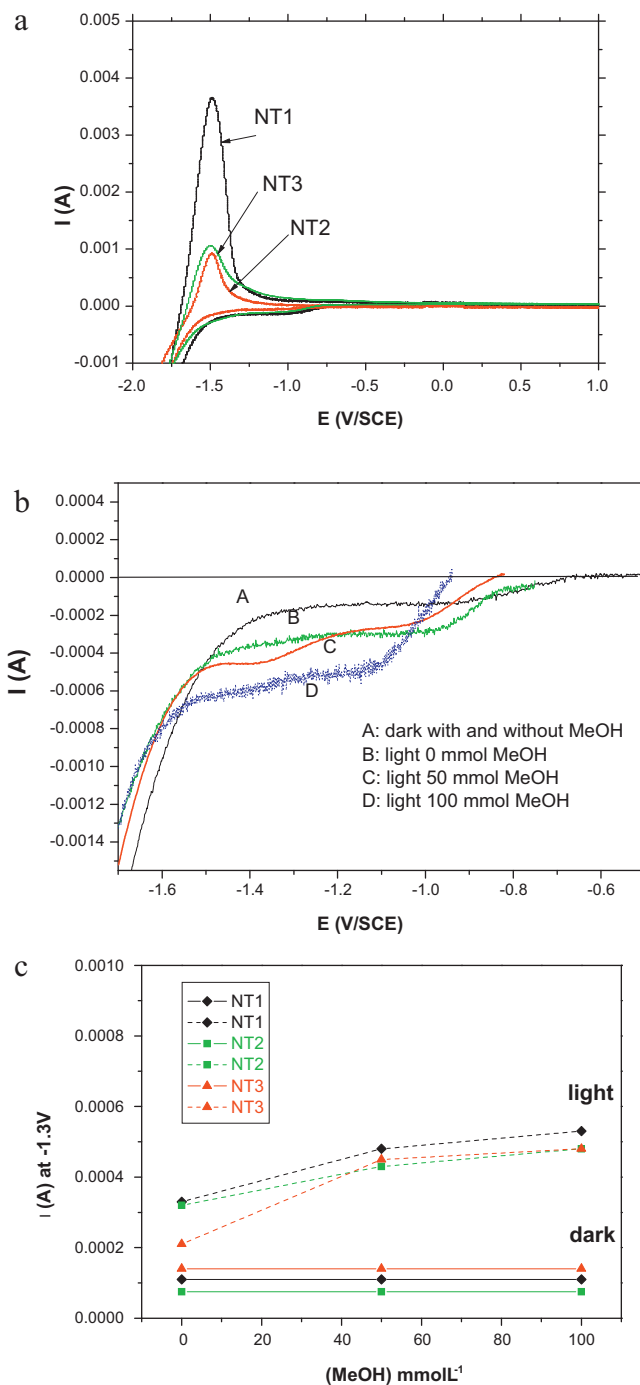
During the reverse scan, a current peak at -1.5 V is observed for the three layers but with different amplitudes.

#### 3.6.1. Cathodic behaviour

We mainly focused on the current plateau due to oxygen reduction. The cathodic reduction of oxygen in the dark has been investigated by the past at several types of compact titanium oxide cathodes. In all cases, it was assumed to occur in alkaline solutions by a 4 electron-reaction into water [40]. Moreover it has been reported to occur through some active sites at the surface. According to Baez et al. [41], oxygen reduction is initiated by the surface reaction:



followed by an electron transfer from Ti(III) to oxygen.



**Fig. 5.** Voltammogram ( $\nu = 50 \text{ mV/s}$ ) in 0.1 M MeOH: (a) for the 3 layers in the cathodic and anodic potential range; (b) for NT<sub>1</sub> in the cathodic potential range in the dark and under lightening; (c) evolution of the cathodic current measured at -1.3 V as a function of [MeOH].

Since oxygen reduction proceeds through the interaction of oxygen with Ti(III), the current for oxygen reduction is likely to be an indication of the amount of Ti(III) sites present in the three types of layers.

Ti(III) sites are produced on a cathodically polarized TiO<sub>2</sub> surface, but have also been shown to be photogenerated under band-gap illumination [30].

As an example, Fig. 5(b) shows the evolution of the current plateau for oxygen reduction in the case of NT1. In the dark, the current value (about  $1.4 \times 10^{-4}$  A) does not depend on the concentration of methanol. On the contrary, under illumination, the current plateau significantly increases and still increases when methanol concentration increases. Similar experiments were performed in the cases of NT<sub>2</sub> and NT<sub>3</sub>, and the results are reported in Fig. 5(c). The 3 layers have similar behaviours.

Since TiO<sub>2</sub> is an n-type semiconductor the photo-induced cathodic current is not a real photocurrent. But if we assume that additional surface states like Ti(III) sites are generated during illumination at open-circuit potential, the increase of the cathodic current is correlated with these additional states. In the dark, the reduction current plateaus for oxygen are very close (75–140  $\mu$ A), and the small differences can be due to slightly different active surfaces. Moreover in the dark, the reduction rate is independent of methanol concentration. Taking into account the geometric surface of the electrodes (0.64 cm<sup>2</sup>), it is noticeable that a current density close to  $2 \times 10^{-4}$  A cm<sup>-2</sup> is also usually measured for compact TiO<sub>2</sub> layers in the same conditions (here not shown). Consequently in the dark, the walls of the nanotubes are not likely to participate in the oxygen reduction reaction, and only the tube bottoms are likely to be active.

*Under light exposure, in absence of methanol:* Under lightening, the current increases for the three layers by a factor of 3, 4 and 2 for respectively NT<sub>1</sub>, NT<sub>2</sub> and NT<sub>3</sub>. If we assume that the current increase is due to activation of the wall surface, NT<sub>1</sub> should show the higher increase due to its higher specific surface (Table 1). Experimentally, this is not the case, since the cathodic current increase is more important for NT<sub>2</sub>, in comparison with NT<sub>1</sub> and NT<sub>3</sub>. This apparent contradiction suggests that during light exposure, more photoinduced surface states for electron storage are generated at the walls of the NT<sub>2</sub> layer, giving rise to new states available for oxygen reduction.

*Under light exposure, in presence of methanol:* The dependence of the cathodic current on methanol concentration is similar for the three types of layers and as previously can be attributed to the hole scavenging role of methanol, which hinders electron–hole recombination at open-circuit potential and allows more photoelectrons to be available for the generation of states at the wall surfaces.

### 3.6.2. Anodic behaviour

The peak appearing at  $-1.5$  V during the reverse potential sweep has often been observed in the literature and has generally been attributed to the reverse of some reactions induced during the cathodic scan, such as oxidation of Ti (III) to the Ti (IV) [38,42] or de-insertion of protons. These two reactions involve energy states in the oxide layer and the surface of the peak is thus an indication of the presence of such states and of the reversibility of the charging–discharging reaction. Quantitative exploitation of this peak can here hardly be done, because the different events are difficult to be separated.

## 3.7. Electrochemical impedance spectrometry

Impedance spectroscopy has been shown to be an effective tool for in situ investigation of charge layer capacitance and accumulation properties in titanium oxide layers [11,12,14,15,23,30,31].

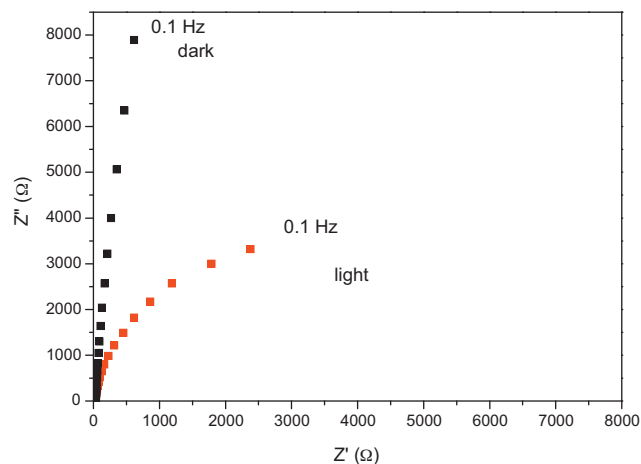


Fig. 6. Impedance diagrams in Nyquist representation for NT<sub>1</sub> nanotubular layer in supporting electrolyte (0.1 mol/L of NaOH) in the dark and under illumination.

Fig. 6 shows impedance diagrams in Nyquist representation for NT<sub>1</sub> samples, in the dark and under illumination, in absence of methanol in the electrolyte. In the dark, the behaviour is mainly capacitive, whereas under illumination an additional step due to charge transfer between TiO<sub>2</sub> and the electrolyte has to be taken into account. A contribution related to the response of energy traps in the nanostructured layer can generate an additional time constant, but its influence on impedance measurements is generally observed in a more anodic potential range [30].

### 3.7.1. In the dark

In the dark where no charge transfer is expected, the behaviour of the oxide layers can be well represented with two RC circuits connected in series [12] or in parallel [14,15]. The first RC circuit accounts for the behaviour of the space charge layer of the semiconducting oxide, whereas the second RC circuit describes relaxation via surface states. In the present work, the equivalent circuit used in order to characterize the electrochemical behaviour of the two layers is reported in Fig. 7(a). A ( $R_{SS}$ ,  $C_{SS}$ ,  $W_{SS}$ ) series circuit corresponding to surface states is added in parallel to a basic  $R_T$ –CPE

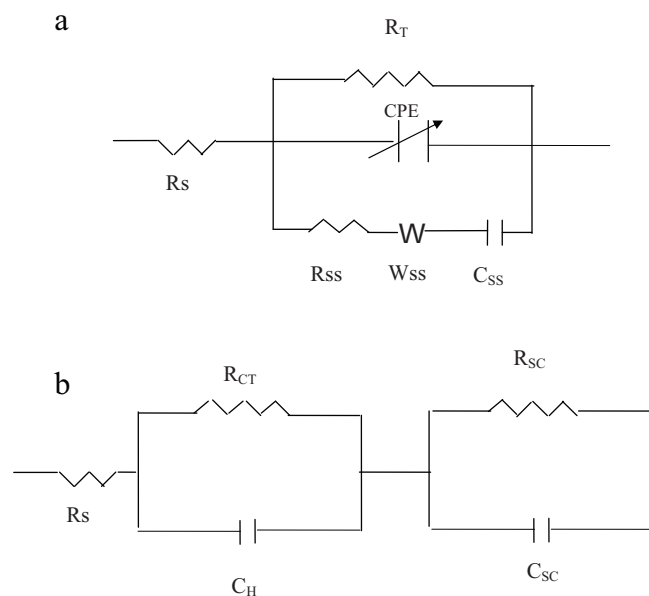


Fig. 7. Equivalent circuits used for impedance diagrams: (a) in the dark; (b) under lightening.

**Table 3**  
Capacitance values  $C_{SC}$  and  $C_{SS}$  extracted from impedance diagrams for the three layers in the dark.

Electrode/dark	[MeOH] mmol L <sup>-1</sup>	$Q_0/\mu(\Omega^{-1} s^n)$	$n$	$C_{SC}/(\mu F)$ calculated from (7)	$C_{SS}/\mu F$
NT <sub>1</sub>	0	47	1	47	145
NT <sub>1</sub>	50	37	1	37	164
NT <sub>1</sub>	100	38	1	38	161
NT <sub>2</sub>	0	21	0.94	20	100
NT <sub>2</sub>	50	39	0.95	28	86
NT <sub>2</sub>	100	35	0.95	25	137
NT <sub>3</sub>	0	34	0.96	25	136
NT <sub>3</sub>	50	23	0.95	16	145
NT <sub>3</sub>	100	38	1	38	154

circuit, and the result in series with an electrolyte resistance  $R_S$ . This model enables data fitting with a relative error for real and imaginary parts of the impedance of less than 1% over the whole frequency range. For the space charge layer of the semiconductor, a constant phase element (CPE) was used instead of a pure capacitance, though the exponent of the CPE was between 1 and 0.95 (Table 3), indicating a near ideal capacitive behaviour.  $R_{SS}$  and  $C_{SS}$  are associated respectively with surface states resistance and capacitance. The Warburg element  $W_{SS}$  is added to obtain a better fit and accounts for some delay in the surface states response. The true charge layer capacitance value  $C_{SC}$  was calculated using the following equation from Brug and co-workers [43]:

$$C = \frac{(RQ_0)^{1/n}}{R} \quad \text{where } R \text{ is defined by } \frac{1}{R} = \frac{1}{R_S} + \frac{1}{R_T} \quad (7)$$

where  $Q_0$  is the constant representative for CPE ( $\Omega^{-1} s^n$ ),  $n$  its exponent.

Impedance measurements were performed at the open-circuit potential for the three layers in the dark, as a function of methanol concentration and the values of the capacitance  $C_{SC}$  and  $C_{SS}$  are reported in Table 3.

It appears that space charge layer capacitance values vary between 16 and 47  $\mu F$  for the three samples and do not depend on methanol concentration. This is in agreement with the conclusion from open-circuit potential measurements in the dark which showed that the values are very close for the three types of layer with and without methanol in the electrolyte (Fig. 3).

Surface state capacitance  $C_{SS}$  between 140 and 160  $\mu F$  for NT<sub>1</sub>, and NT<sub>3</sub>, is slightly lower (between 80 and 140  $\mu F$ ) for NT<sub>2</sub> but keeps the same order of magnitude for the 3 methanol concentrations. It is noticeable that a Randle circuit with a simple  $R_p$ -CPE circuit in series with the electrolyte resistance can also be used, yet with a slightly higher fitting error. In the latter case, the equivalent capacitance is close to the sum of the parallel capacitances  $C_{SC}$  and  $C_{SS}$ .

### 3.7.2. Under illumination

As shown in Fig. 3, in passing from dark to illumination the open-circuit potential reaches a very negative value which is then close to the conduction band. The equivalent circuit used in this case is reported in Fig. 7(b). A ( $R_{CT}$ ,  $C_H$ ) circuit corresponding to charge transfer ( $R_{CT}$  is the charge transfer resistance and  $C_H$  is the capacitance of the Helmholtz layer) is added in series with a  $R_{SC}$ - $C_{SC}$  circuit, and the result in series with an electrolyte resistance  $R_S$ . This model enables data fitting with a relative error for real and imaginary parts of the impedance of less than 2% over the whole frequency range. For the semiconductor, the capacitance  $C_{SC}$  includes the space charge layer and the surface states capacitance. Since in the potential range close to the conduction band most energy states in the gap are filled with photogenerated electrons, the  $C_{SC}$  value can be considered to be mainly due to the space charge layer.  $C_H$ ,  $C_{SC}$  and  $R_{CT}$  values during light exposure are reported in Table 4 as a

function of methanol concentration, for the 3 types of nanotubular layers.

The  $R_{CT}$  values have the same order of magnitude for the three types of layers. The decrease of  $R_{CT}$  when methanol concentration increases is correlated with an increase of the rate of charge transfer as expected.

In passing from dark to illumination an increase of  $C_{SC}$  is observed which is similar for NT<sub>1</sub>, NT<sub>2</sub> and NT<sub>3</sub>. This is partly due to the shift of the open-circuit potential towards more cathodic, but can also be attributed to the increase of the active surface since UV excitation of the nanotubes induces activation of the tube walls [30]. A noticeable result is the increase of the semiconductor capacitance with methanol concentration, which can also be correlated to the higher photopotential in presence of higher concentrations of MeOH (Fig. 3). This effect was attributed above to the fact that a higher methanol concentration and consequently a higher amount of adsorbed methanol at the surface of the tubes allows more electrons to accumulate in the electrode, generating Ti(III) species as surface states.

Another noticeable result is the huge value of the double layer capacitance, which, in aqueous solutions, is usually close to 50  $\mu F/cm^2$ . That clearly shows that the whole surface, including the tube walls, is active. In Table 4, if the extracted  $C_H$  values are divided by the specific surface reported in Table 1, values of double layer capacitance between 28 and 85  $\mu F/cm^2$  are determined. This is an additional proof that the whole surface of the tubes is activated during lightening.

The  $C_{SC}$  values divided by the specific surface of each layer are also reported in Table 4. It appears that NT<sub>2</sub> shows the higher capacitance density, yet the lower photopotential (Fig. 3). As stated previously, if we assume that  $C_{SC}$  involves simultaneously the charge layer capacitance and the surface state capacitance, the high value under illumination of  $C_{SC}$  in the case of NT<sub>2</sub> is likely to be due to the generation of a higher amount of surface states in comparison with NT<sub>1</sub> and NT<sub>3</sub>. This is in agreement with the observations made during cathodic reduction of oxygen via surface states (see Section 3.6.1).

### 3.8. Photocatalytic (PC) degradation of AB

In order to compare the photocatalytic efficiency of the three nanotubular TiO<sub>2</sub> electrodes, the variation of relative concentration of Amino Black (AB) as a function of reaction time was carried out and shown in Fig. 8(a).

It was observed that the degradation of dye proceeded much faster in the presence of NT<sub>1</sub> as compared with the two others TiO<sub>2</sub> nanotubes samples. Fig. 8(b) shows that a 96.6% decolouration efficiency of dye in aqueous solution was achieved after 45 min of reaction using NT<sub>1</sub>, while the decolourization efficiencies obtained for NT<sub>2</sub> and NT<sub>3</sub> were only 73.5 and 80.0%, respectively, within the same time. It is noticeable that the degradation rate measured on a TiO<sub>2</sub> compact layer (not shown here) in the same experimental

**Table 4**  
Capacitance  $C_H$  and  $C_{SC}$ , and charge transfer resistance values extracted from impedance diagrams for the three layers under lightening.

Electrode/light exposure	[MeOH] (mmol L <sup>-1</sup> )	$C_H$ ( $\mu$ F)	$C_H/S_{spec}$ ( $\mu$ F/cm <sup>2</sup> )	$R_{CT}$ ( $\Omega$ )	$C_{SC}$ ( $\mu$ F)	$C_{SC}/S_{spec}$ ( $\mu$ F/cm <sup>2</sup> )
NT <sub>1</sub>	0	1970	50	20	315	8.1
NT <sub>1</sub>	50	1973	50	18	325	8.3
NT <sub>1</sub>	100	1104	28	4	676	17.3
NT <sub>2</sub>	0	1961	85	13	301	13
NT <sub>2</sub>	50	1160	50	11	330	14.3
NT <sub>2</sub>	100	1078	47	4	591	26.7
NT <sub>3</sub>	0	1667	57	37	259	8.9
NT <sub>3</sub>	50	883	30	5	380	13.1
NT <sub>3</sub>	100	1083	36	4	620	21.3

conditions lead to an efficiency of only 17% after 45 min of treatment.

The reason often put forward to explain the good photocatalytic behaviour of NT layers is the large surface area, which provides high contact area for absorption of dye or pollutants. In the present work, the solid fraction depending on the voids between the tubes and the thickness of the tube walls was also taken into account to estimate the level of excitation in the oxide layer. Photocurrent measurements in the saturation range clearly showed that the difference in efficiency between the three types of layers can be attributed to their differences in morphology: specific surface but also solid fraction available for generated charges (Tables 1 and 2). Since the photocatalytic performances are correlated with the photocurrent,

the comparison between NT<sub>1</sub>, NT<sub>2</sub>, NT<sub>3</sub> leads to attribute the higher photocatalytic efficiency in NT<sub>1</sub> to a higher ( $S_{spec} S_{\%} L\alpha$ ) product under UV light. Since the tube wall thickness is similar for the three types of layers, the influence of this parameter on the photocatalytic performances cannot be confirmed in the present work, as was suggested in Ref. [28].

In passing from dark to light exposure, potentiodynamic measurements in the cathodic potential range evidenced the higher increase in the rate of oxygen reduction for NT<sub>2</sub> in comparison with NT<sub>1</sub> and NT<sub>3</sub>, due to a higher amount of photo-generated surface states. Accordingly, impedance measurements performed at open-circuit potential showed a higher capacitance increase for the NT<sub>2</sub> layer under light exposure. From these results it seems that the photo-catalytic performances of a NT layer are mainly determined by the specific surface and the solid fraction of the layer in which photo-charges are created. They are not influenced by the photo-generated surface states in which photoelectrons can be stored.

#### 4. Conclusion

The nanotubular TiO<sub>2</sub> electrodes were obtained by anodization of a Ti sheet at 20 V for 45 min in three different media. The three types of layers showed similar thickness (between 800 and 1700 nm) and the tube-walls had the same thickness (between 15 and 20 nm). The titania nanotube electrode prepared in ethylene glycol (NT<sub>1</sub>) showed a higher photocurrent, and a higher activity for pollutant degradation. This result is attributed to a good compromise between a high specific surface and the higher solid fraction able to absorb incident photons to convert into hole–electron pairs.

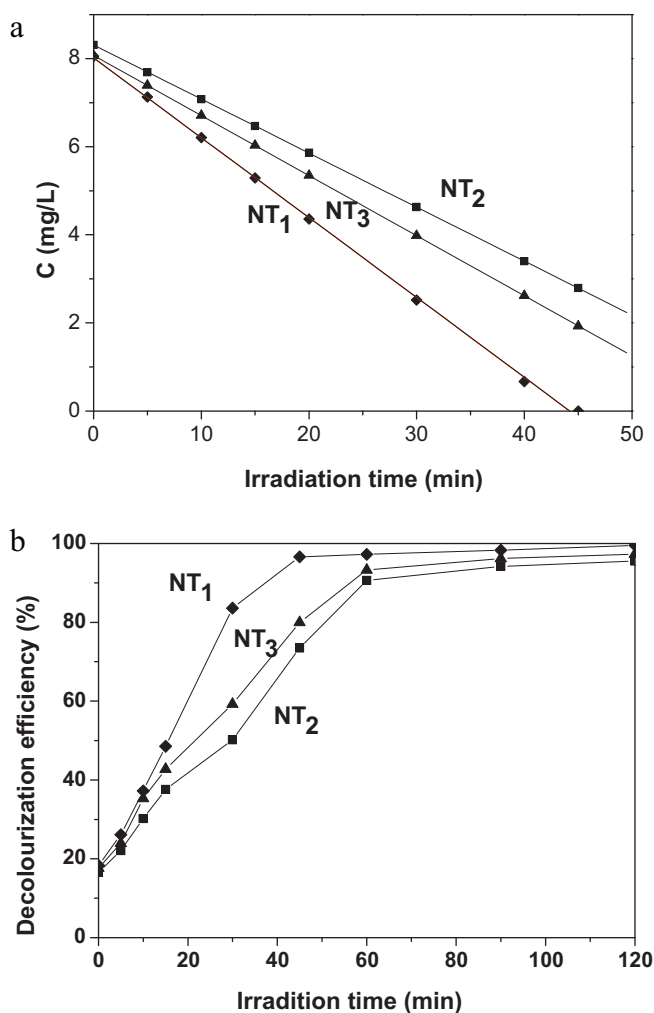
If we consider the amount of photo-generated traps in the tube walls, it appears that in passing from dark to UV exposure, the NT<sub>2</sub> layer characterized by rough walls, exhibits the higher increase in surface state generation, and this result is also confirmed through impedance measurements, by a higher capacitance density increase under light exposure. From these first results, it seems that photo-generated surface states have no influence on the photocatalytic activity of a nanotubular layer. The main reason is that photo-generated surface states able to trap photoelectrons are different from the states at which holes can recombine or through which they are transferred to the electrolyte. But further experiments are required to confirm this assumption.

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**Fig. 8.** (a) PC degradation of black amido ( $C_0 = 10$  mg/L,  $\lambda = 618$   $\mu$ m); (b) decolourization efficiency of Black Amido.



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